



# The $\rho$ -Loewner Energy

Large deviations, minimizers, and alternative descriptions

Ellen Krusell, KTH Royal Institute of Technology

Supervisor: Fredrik Viklund

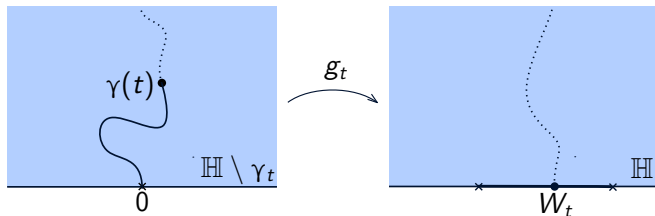
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# The Chordal Loewner Energy

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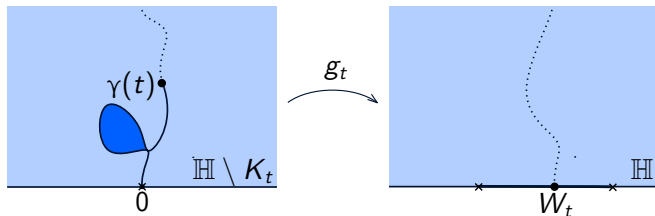
# Chordal Loewner Evolution



- Chord  $\gamma$  in  $(\mathbb{H}; 0, \infty)$ . Denote  $\gamma_t := \gamma([0, t])$ .
- **Mapping-out functions**  $g_t : \mathbb{H} \setminus \gamma_t \rightarrow \mathbb{H}$  satisfying,  
 $g_t(z) = z + \frac{a(t)}{z} + o(\frac{1}{z})$  as  $z \rightarrow \infty$ . Parametrize  $\gamma$  s.t.  $a(t) = 2t$ .
- The **driving function**,  $W_t := g_t(\gamma(t))$ , encodes  $\gamma$  via the chordal **Loewner differential equation**:

$$\partial_t g_t(z) = \frac{2}{g_t(z) - W_t}, \quad g_0(z) = z, \quad (1)$$

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- Solving (1) for an arbitrary continuous function  $t \mapsto W_t$  gives family of locally growing compact sets  $(K_t)_{t \geq 0}$ .

# Schramm-Loewner Evolution

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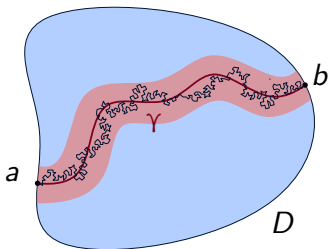
- Chordal **Schramm-Loewner evolution**,  $SLE_\kappa$ , in  $(\mathbb{H}; 0, \infty)$  is the random curve driven by  $W_t = \sqrt{\kappa}B_t$ ,  $\kappa \geq 0$ , where  $B_t$  standard Brownian motion [Schramm '00]. A.s. simple when  $\kappa \leq 4$ .
- Chordal SLE in  $(D; a, b)$  is defined by conformal mapping  $(D; a, b) \rightarrow (\mathbb{H}; 0, \infty)$ .
- $SLE_\kappa$  is the scaling limit of interfaces in planar critical lattice models for several values of  $\kappa > 0$ .
- $SLE_0$  is the hyperbolic geodesic in  $(D; a, b)$ . ( $W \equiv 0 \rightsquigarrow \gamma = i\mathbb{R}^+$ .)

# Schramm-Loewner Evolution

Chordal  $\text{SLE}_\kappa$  satisfies the large deviation principle as  $\kappa \rightarrow 0+$

$$\mathbb{P}[\text{SLE}_\kappa \text{ stays close to } \gamma] \approx \exp\left(-\frac{I^{(D;a,b)}(\gamma)}{\kappa}\right), \text{ as } \kappa \rightarrow 0+$$

where  $I^{(D;a,b)}$  is the **chordal Loewner energy**.



[Wang '19, Peltola-Wang '23, Guskov '23, Abuzaid-Peltola '24]

# Chordal Loewner Energy

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The **chordal Loewner energy** of a chord  $\gamma$  in  $(\mathbb{H}; 0, \infty)$  is defined by

$$I^{(\mathbb{H}; 0, \infty)}(\gamma) = \frac{1}{2} \int_0^\infty (\dot{W}_t)^2 dt,$$

if  $W$  is absolutely continuous, and  $I^{(\mathbb{H}; 0, \infty)}(\gamma) = \infty$  otherwise [Wang '19].

- The unique minimizer of  $I^{(\mathbb{H}; 0, \infty)}$  is  $\text{SLE}_0$ , that is  $i\mathbb{R}^+$ .
- $I^{(D; a, b)}$  is defined on chords in  $(D; a, b)$  by conformal mapping  $(D; a, b) \rightarrow (\mathbb{H}; 0, \infty)$ .
- Other description(s)?

# Dirichlet Energy Formula

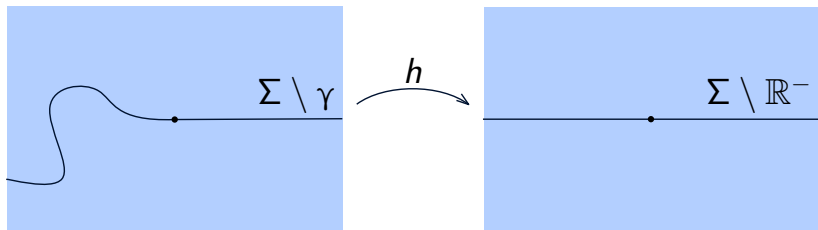
Let  $\Sigma = \mathbb{C} \setminus \mathbb{R}^+$  and consider a chord  $\gamma$  in  $(\Sigma; 0, \infty)$ .

## Theorem (Wang '19)

*The chordal Loewner energy can be expressed as,*

$$I^{(\Sigma; 0, \infty)}(\gamma) = \frac{1}{\pi} \int_{\Sigma \setminus \gamma} |\nabla \log |h'||^2 dz^2,$$

where  $h : \Sigma \setminus \gamma \rightarrow \Sigma \setminus \mathbb{R}^-$  is conformal and  $h(\infty) = \infty$ .

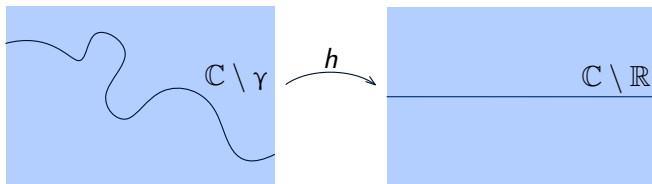




## Loop Loewner Energy

The **loop Loewner energy** of a Jordan curve  $\gamma$  rooted at  $\infty$  on  $\hat{\mathbb{C}}$  is

$$I^L(\gamma) = \frac{1}{\pi} \int_{\hat{\mathbb{C}} \setminus \gamma} |\nabla \log |h'||^2 dz^2.$$



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- $I^L$  is invariant of the root [Rohde-Wang '21].
- Link to **Teichmüller theory**:  $I^L$  coincides (up to a multiplicative constant) with the **universal Liouville action** [Wang '19].
- The class of finite Loewner energy loops = The class of **Weil-Petersson quasicircles** [Wang '19].
- Loewner energy also shows up in  $n \rightarrow \infty$  limit of partition functions of **Coulomb gas** on a Jordan curve/domain [Johansson '22, Johansson-Viklund '23].

**SLE<sub>κ</sub>(ρ)**

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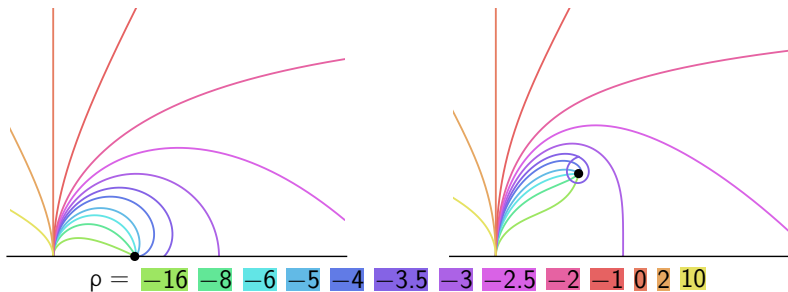
# SLE<sub>κ</sub>(ρ)

- **Chordal SLE<sub>κ</sub>(ρ)**,  $\rho \in \mathbb{R}$ , in  $(\mathbb{H}; 0, \infty)$ , with force point  $z_0 \in \overline{\mathbb{H}} \setminus \{0\}$ , is the random curve whose driving function satisfies the SDE:

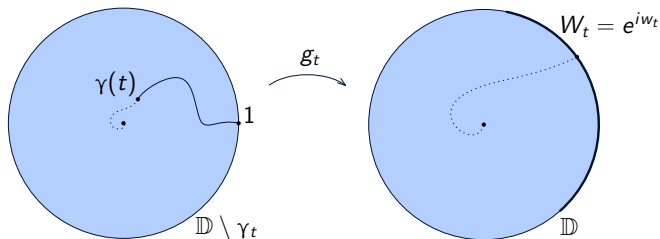
$$dW_t = \operatorname{Re} \frac{\rho}{W_t - z_t} dt + \sqrt{\kappa} dB_t, \quad dz_t = \frac{2}{z_t - W_t} dt,$$

where  $z_t = g_t(z_0)$ . Defined up to  $\tau = \lim_{\epsilon \rightarrow 0} \inf\{t : |W_t - z_t| \leq \epsilon\}$ .

- $\kappa = 0$ :



# Radial Loewner Evolution



- Slit  $\gamma$  in  $(\mathbb{D}; 1, 0)$ , parametrized by conformal radius.
- **Mapping-out functions**  $g_t : \mathbb{D} \setminus \gamma_t \rightarrow \mathbb{D}$ ,  $g_t(0) = 0$ ,  $g_t'(0) = e^t$ .
- Radial **Loewner differential equation**,

$$\partial_t g_t(z) = g_t(z) \frac{W_t + g_t(z)}{W_t - g_t(z)}, \quad g_0(z) = z,$$

where  $W_t = e^{i w_t} = g_t(\gamma(t))$ .

- $w_t = \sqrt{\kappa} B_t \rightsquigarrow$  radial  $\text{SLE}_\kappa$  in  $(\mathbb{D}; 1, 0)$ .

- **Radial SLE <sub>$\kappa$</sub> ( $\rho$ )**, in  $(\mathbb{D}; 1, 0)$ , with force point  $z_0 = e^{iv_0}$ ,  $v_0 \in (0, 2\pi)$  is the random curve whose driving function  $W_t = e^{iw_t}$  satisfies the SDE:

$$dw_t = \frac{\rho}{2} \cot\left(\frac{w_t - v_t}{2}\right) dt + \sqrt{\kappa} dB_t, \quad dv_t = \cot\left(\frac{v_t - w_t}{2}\right) dt,$$

where  $e^{iv_t} = g_t(e^{iv_0})$ .

- First appearance: SLE<sub>8/3</sub>( $\rho$ ),  $\rho > -2$ , is the outer boundary of sets satisfying a conformal restriction property [Lawler-Schramm-Werner '03].
- Main player in imaginary geometry: generalized flow-line of GFF [Miller-Sheffield '16, '17].
- Coordinate change property: An SLE <sub>$\kappa$</sub> ( $\rho$ ) in  $(D; a, b)$  with force point  $c$ , is an SLE <sub>$\kappa$</sub> ( $\kappa - 6 - \rho$ ) in  $(D; a, c)$  with force point  $b$  [Schramm-Wilson '05].

# The $\rho$ -Loewner Energy

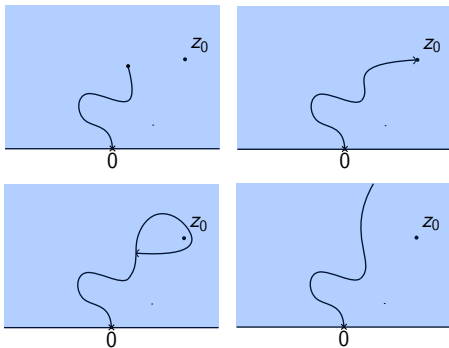
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# Chordal $\rho$ -Loewner Energy

- Fix  $\rho \in \mathbb{R}$  and  $z_0 \in \overline{\mathbb{H}} \setminus \{0\}$ . The **chordal  $\rho$ -Loewner energy** of a simple curve  $\gamma$  starting at 0 in  $\mathbb{H} \setminus \{z_0\}$  is defined as

$$I_{\rho, z_0}^{(\mathbb{H}; 0, \infty)}(\gamma) = \frac{1}{2} \int_0^T \left( \dot{W}_t - \operatorname{Re} \frac{\rho}{W_t - z_t} \right)^2 dt,$$

if  $W_t$  is absolutely continuous, and  $I_{\rho, z_0}^{(\mathbb{H}; 0, \infty)}(\gamma) = \infty$  otherwise.





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if  $W_t$  is absolutely continuous, and  $I_{\rho, z_0}^{(\mathbb{H}; 0, \infty)}(\gamma) = \infty$  otherwise.

- Consistent with Freidlin–Wentzell theory  
( $dW_t = \operatorname{Re} \frac{\rho}{W_t - z_t} dt + \sqrt{\kappa} dB_t$ ).
- If  $z_0 \in \mathbb{H}$  and  $T < \tau$  (i.e.  $\inf_{t \leq T} |W_t - z_t| > 0$ )

$$I_{\rho, z_0}^{(\mathbb{H}; 0, \infty)}(\gamma) = I^{(\mathbb{H}; 0, \infty)}(\gamma) + \rho \log \frac{\sin \theta_T}{\sin \theta_0} - \frac{\rho(8 + \rho)}{8} \log \frac{|g'_T(z_0)| y_T}{y_0},$$

where  $y_T = \operatorname{Im}(z_T)$  and  $\theta_T = \arg(z_T)$ .

- If  $T < \tau$ , then  $I_{\rho, z_0}^{(\mathbb{H}; 0, \infty)}(\gamma) < \infty$  iff  $I^{(\mathbb{H}; 0, \infty)}(\gamma) < \infty$ .
- Similar formula when  $z_0 \in \mathbb{R} \setminus \{0\}$ .

## Radial $\rho$ -Loewner energy

- Fix  $\rho \in \mathbb{R}$  and  $v_0 \in (0, 2\pi)$ . The **radial  $\rho$ -Loewner energy** of a simple curve starting at 1 in  $\mathbb{D} \setminus \{0\}$  is defined as

$$I_{\rho, e^{iv_0}}^{(\mathbb{D}; 1, 0)}(\gamma) = \frac{1}{2} \int_0^T \left( \dot{w}_t - \frac{\rho}{2} \cot \left( \frac{w_t - v_t}{2} \right) \right)^2 dt,$$

if  $w_t$  is absolutely continuous, and  $I_{\rho, e^{iv_0}}^{(\mathbb{D}; 1, 0)}(\gamma) = \infty$  otherwise.

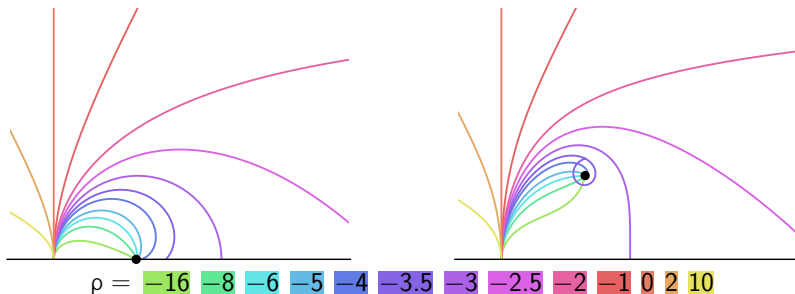
- $I_{\rho, c}^{(D; a, b)}$  is defined via conformal map to  $\mathbb{H}$  or  $\mathbb{D}$ .
- Coordinate change property:

$$I_{\rho, c}^{(D; a, b)}(\gamma) = I_{-6-\rho, b}^{(D; a, c)}(\gamma).$$

- If  $T < \tau$ , then  $I_{\rho, e^{iv_0}}^{(\mathbb{D}; 1, 0)}(\gamma) < \infty$  iff  $I_{\rho, e^{iv_0}}^{(\mathbb{D}; 1, e^{iv_0})}(\gamma) < \infty$ .

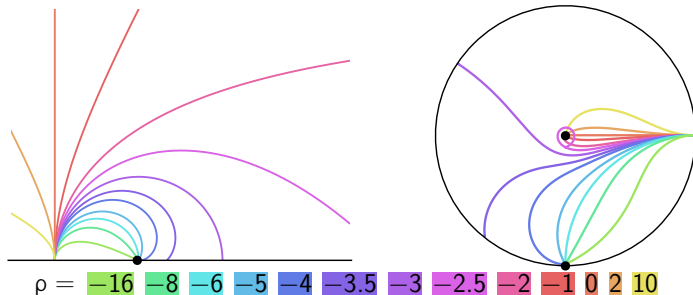
# Minimizers

The unique minimizer of the  $\rho$ -Loewner energy is the  $SLE_0(\rho)$  curve.



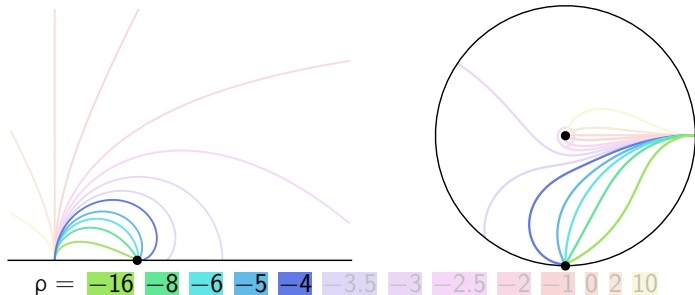
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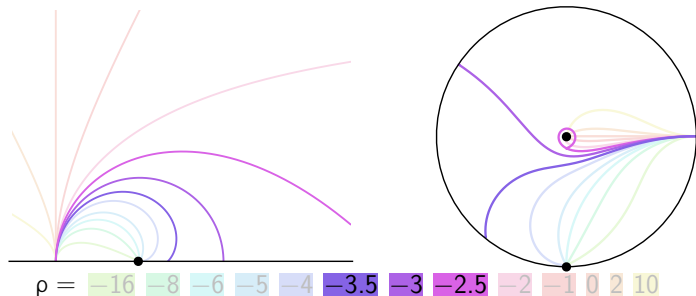


Three phases:

- When  $\rho \in (-\infty, -4]$  :  $SLE_0(\rho)$  hits the force point.

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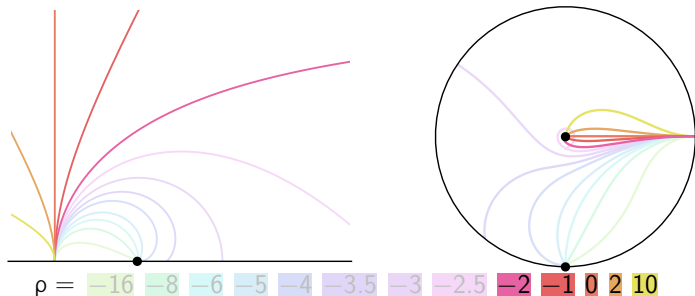


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- When  $\rho \in (-\infty, -4]$ :  $SLE_0(\rho)$  hits the force point.
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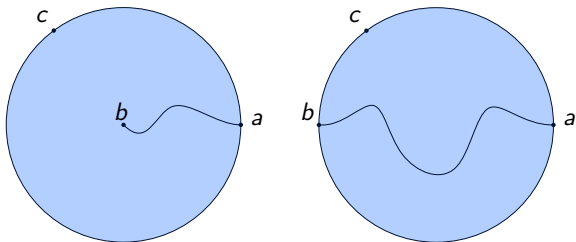
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- When  $\rho \in (-4, -2)$ :  $SLE_0(\rho)$  separates the reference point and force point.
- When  $\rho \in [-2, \infty)$ :  $SLE_0(\rho)$  approaches the reference point.

# Large Deviation Principle

Fix  $\rho > -2$ ,  $a \in \partial\mathbb{D}$ ,  $b \in \bar{\mathbb{D}} \setminus \{a\}$ , and  $c \in \partial\mathbb{D} \setminus \{a, b\}$ .

## Theorem (K '24)

The  $SLE_\kappa(\rho)$  in  $(\mathbb{D}; a, b)$  with force point  $c$  satisfies a large deviation principle as  $\kappa \rightarrow 0+$ , with respect to the Hausdorff topology on the space of simple curves  $\gamma$  in  $(\mathbb{D}; a, b)$ , with good rate function  $I_{\rho,c}^{(\mathbb{D};a,b)}$ .



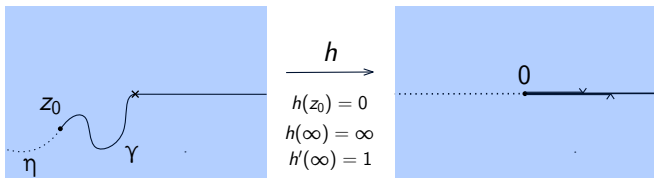
- Proof follows the same outline as in [Peltola-Wang '23]: LDP on driving process on  $[0, T] \rightsquigarrow$  LDP on finite time curves  $\rightsquigarrow$  LDP on infinite time curves.



# Dirichlet Energy Formulas

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# Radial Case



## Lemma (K '24)

If  $I^{(\Sigma;0,\infty)}(\gamma) := I^{(\Sigma;0,\infty)}(\gamma \cup \eta) < \infty$ , then

$$|h'(z_0)|_{\eta} := \lim_{\epsilon \rightarrow 0^+} \frac{|h(\eta(\epsilon)) - h(z_0)|}{|\eta(\epsilon) - z_0|}$$

exists and is positive.

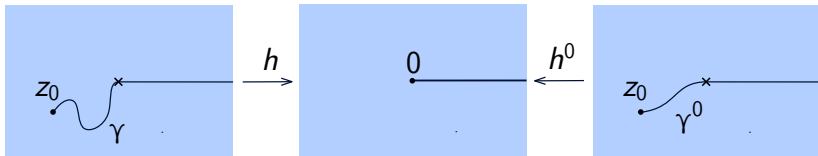
# Radial Case

## Theorem (K '24)

Fix  $\rho > -2$ . Let  $\gamma^0$  denote the  $SLE_0(\rho)$  in  $(\Sigma; 0, z_0)$  with force point at  $\infty$ , and let  $h^0$  be the corresponding conformal map. A simple curve  $\gamma$  in  $(\Sigma; 0, z_0)$  has finite  $\rho$ -Loewner energy w.r.t.  $\infty$  if and only if  $\mathcal{D}(\log |h'|) < \infty$ , in which case,

$$I_{\rho, \infty}^{\Sigma; 0, z_0}(\gamma) = \mathcal{D}(\log |h'|) - \mathcal{D}(\log |(h^0)'|) - \frac{(\rho + 6)(\rho - 2)}{8} \log |H'(z_0)|_{\eta}$$

where  $H = (h^0)^{-1} \circ h$ , and  $\mathcal{D}(f) = \frac{1}{\pi} \int |\nabla f|^2 dz^2$ .



## Radial Case

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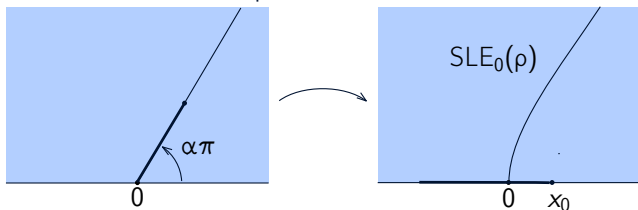
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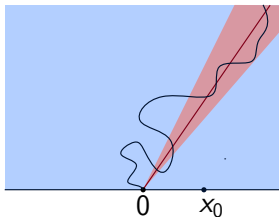
Proof uses coordinate change property, integrated formula, and the Dirichlet energy formula for the chordal Loewner energy.

## Chordal Case

- $\text{SLE}_0(\rho)$ ,  $\rho > -2$ , with force point  $x_0 > 0$ , is a “mapped out ray” of angle  $\alpha\pi$ , where  $\alpha = \alpha(\rho) = \frac{\rho+2}{\rho+4} \in (0, 1)$ .



- Curves of finite  $\rho$ -Loewner energy,  $\rho > -2$ , approach  $\infty$  within cones around angle  $\alpha\pi$



## Chordal Case

For a chord  $\gamma$  in  $(\Sigma; 0, \infty)$ , let  $h : \Sigma \setminus \gamma \rightarrow \Sigma \setminus \mathbb{R}^-$  be conformal with  $h(\infty) = \infty$ . We define, whenever the limit exists,

$$\mathcal{D}_\beta(\log |h'|) = \lim_{R \rightarrow \infty} \left( \frac{1}{\pi} \int_{B(0,R)} |\nabla \log |h''||^2 dz^2 - c_\beta \log R \right), \quad c_\beta = \frac{(1 - 2\beta)^2}{2\beta(1 - \beta)}.$$

### Theorem (K '24)

Fix  $\rho > -2$  and  $x_0 > 0$ . If  $\gamma$  is a chord in  $(\Sigma; 0, \infty)$ , such that there exists  $T$  for which  $\gamma([T, \infty))$  is the  $\rho$ -Loewner energy optimal extension of  $\gamma_T$ , then

$$I_{\rho, x_0^+}^{(\Sigma; 0, \infty)}(\gamma) = \mathcal{D}_\alpha(\log |h'|) - \mathcal{D}_\alpha(\log |(h^0)'|) - \frac{\rho(\rho + 4)}{4} \log |H'(x_0)|,$$

where  $H$  is the conformal map from the upper component of  $\Sigma \setminus \gamma$  to the upper component of  $\Sigma \setminus \gamma^0$  fixing  $\infty$  and  $x_0$  and satisfying  $H'(\infty) = 1$ .

# Determinants of Laplacians

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# Determinants of Laplacians

On smooth compact Riemannian surfaces with boundary,  $(M, g)$ , we may define the  $\zeta$ -regularized determinant of the Dirichlet Laplace-Beltrami operator,  $\det_{\zeta} \Delta_{(M,g)}$  [Ray-Singer '71]:

- If  $\partial M \neq \emptyset$ , then  $0 < \lambda_1 \leq \lambda_2 \leq \dots$ ,  $\lim_{n \rightarrow \infty} \lambda_n \rightarrow \infty$ .
- Define  $\zeta(s) = \sum_{n=1}^{\infty} \lambda_n^{-s}$ , for  $\text{Res} > 1$ . Analytic cont. to nbhd of 0.
- We define  $\det_{\zeta} \Delta_{(M,g)} := e^{-\zeta'(0)}$  justified by

$$\zeta'(s) = - \sum_{n=1}^{\infty} \log \lambda_n \lambda_n^{-s}, \quad \text{Res} > 1 \quad \rightsquigarrow \quad \zeta'(0) = - \log \left( \prod_{n=1}^{\infty} \lambda_n \right).$$

- $\det_{\zeta} \Delta_{(M,g)}$  can be defined in the same way for curvilinear polygonal domains (boundary is p.w. smooth with corner angles in  $(0, \infty)$ ) [Aldana-Kirsten-Rowlett '20, K '24].



# Determinants of Laplacians

Let  $K$  be a compact subset of  $M$  and set (whenever the RHS is defined)

$$\mathcal{H}_{(M,g)}(K) := \log \det_{\zeta} \Delta_{(M,g)} - \sum_{C_i} \log \det_{\zeta} \Delta_{(C_i,g)}$$

where the sum runs over connected components of  $M \setminus K$ .

## Theorem (Wang '19)

*Let  $\gamma$  be a  $C^\infty$  Jordan curve on  $S^2$ , and let  $g$  be a metric conformally equivalent to the spherical metric  $g_0$ , that is,  $g = e^{2\sigma} g_0$ ,  $\sigma \in C^\infty(S^2)$ .*

*Then,*

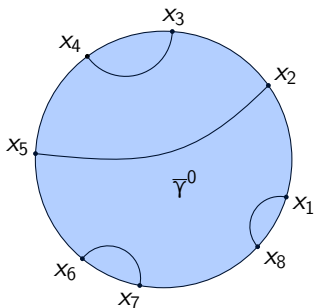
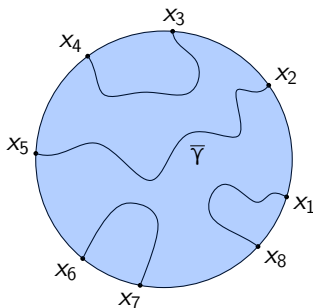
$$I^L(\gamma) = 12(\mathcal{H}_{(S^2,g)}(\gamma) - \mathcal{H}_{(S^2,g)}(S^1)).$$

# Determinants of Laplacians

In [Peltola-Wang '23] it is proved that the **multichordal Loewner energy** can be expressed as

$$I_{\mathbb{D}; x_1, \dots, x_{2n}}^{\alpha}(\bar{\gamma}) = 12(\mathcal{H}_{(\mathbb{D}, g)}(\bar{\gamma}) - \mathcal{H}_{(\mathbb{D}, g)}(\bar{\gamma}^0)),$$

for a smooth and finite energy multichord  $\bar{\gamma}$  with link-pattern  $\alpha$  and  $g = e^{2\sigma} dz^2$ ,  $\sigma \in C^\infty(\bar{\mathbb{D}})$ , where  $\bar{\gamma}^0$  is the geodesic multichord with the same link-pattern.

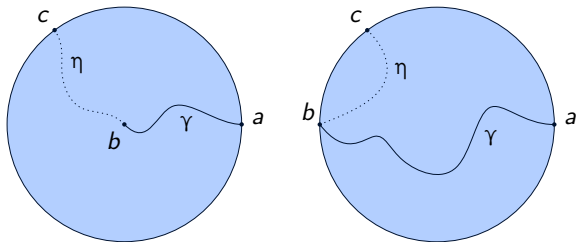


# Determinants of Laplacians

Fix  $\rho > -2$ ,  $a, c \in \partial\mathbb{D}$  and  $b \in \overline{\mathbb{D}}$ . Let  $\gamma$  be a curve in  $(\mathbb{D}; a, b)$  and let  $\eta = \eta(\gamma)$  denote the hyperbolic geodesic from  $b$  to  $c$  in  $\mathbb{D} \setminus \gamma$ . The  $\rho$ -**Loewner potential** of  $\gamma$  with respect to a metric  $g = e^{2\sigma} dz^2$ , is defined as

$$\mathcal{H}_{\rho,c}^{(\mathbb{D};a,b)}(\gamma; g) := \beta \mathcal{H}_{(\mathbb{D},g)}(\gamma) + (1 - \beta) \mathcal{H}_{(\mathbb{D},g)}(\gamma \cup \eta),$$

where  $\beta = \frac{(2-\rho)(\rho+6)}{12}$ , whenever the RHS is well-defined.



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- If  $\rho = 0$ , then  $\beta = 1$ .
- If  $\rho = 2$ , then  $\beta = 0$ . ( $\rho = 2$  corresponds to one arm of a 2-radial or 2-chordal, and  $\eta$  is the “optimal second arm”.)

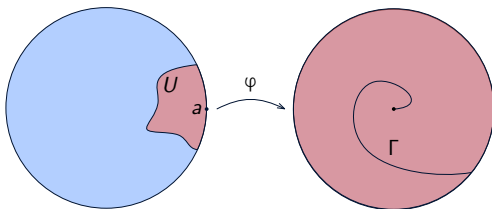
# Determinants of Laplacians

## Proposition (K '24)

For all  $g \in \mathcal{G}(a)$  and  $\gamma_1, \gamma_2 \in \mathcal{X}(g, a, b, c, \rho)$  we have that

$$I_{\rho,c}^{(\mathbb{D};a,b)}(\gamma_1) - I_{\rho,c}^{(\mathbb{D};a,b)}(\gamma_2) = 12(\mathcal{H}_{\rho,c}^{(\mathbb{D};a,b)}(\gamma_1; g) - \mathcal{H}_{\rho,c}^{(\mathbb{D};a,b)}(\gamma_2; g)).$$

- $\mathcal{G}(a)$  consists of metrics  $g = e^{2\sigma} dz^2$ ,  $\sigma \in C^\infty(\overline{\mathbb{D}} \setminus \{a\})$  for which there is a coordinate  $\varphi : U \rightarrow \mathbb{D} \setminus \Gamma$  in a nbhd of  $a$ , for which  $\varphi^*g = e^{2\tilde{\sigma}(w)} dw^2$  and  $\tilde{\sigma}$  extends smoothly to both sides of the slit  $\Gamma$ .



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- $\mathcal{X}(g, a, b, c, \rho)$  consists of  $C^\infty$ -smooth curves  $\gamma$  in  $(\mathbb{D}; a, b)$  which are smoothly attached at  $a$ , and
  - if  $b \in \mathbb{D}$ , then  $\gamma \cup \eta$  is  $C^\infty$ -smooth.
  - if  $b \in \partial\mathbb{D}$ , then there exists  $T$  s.t.  $\gamma([T, \infty))$  is the  $\rho$ -optimal extension of  $\gamma_T$ .
- If  $b \in \partial\mathbb{D}$  or  $b = 0$  and  $a$  and  $c$  are antipodal, then there exists  $g \in \mathcal{G}(a)$  so that the minimizer  $\gamma^0$  belongs to  $\mathcal{X}(g, a, b, c, \rho)$ .

**Thank you for your attention!**

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