# From Partial and Horizontal Contraction to $k$-Contraction 

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#### Abstract

A geometric generalization of contraction theory called $k$-contraction was recently developed using $k$-compound matrices. In this note, we focus on the relations between $k$ contraction and two other generalized contraction frameworks: partial contraction (also known as virtual contraction) and horizontal contraction. We show that in general these three notions of contraction are different. We here provide new sufficient conditions guaranteeing that partial contraction implies horizontal contraction, and that horizontal contraction implies $k$ contraction. We use the Andronov-Hopf oscillator to demonstrate some of the theoretical results.


Index Terms- $k$-contraction, compound matrix, partial contraction, horizontal contraction, virtual contraction, AndronovHopf oscillator

## I. Introduction

Contraction theory is a powerful tool for analyzing the asymptotic behavior of nonlinear time-varying dynamical systems [13], [2], [10]. A contractive system behaves in many respects like a uniformly asymptotically stable linear system: initial conditions are "forgetten" and any two trajectories approach each other at an exponential rate.
There exist easy to verify sufficient conditions for contraction that are based on matrix measures [2] and contraction analysis has found numerous applications such as control synthesis for regulation [18] and tracking [32], observer design [14], [22], [1], optimization [29], synchronization of multi-agents systems [24], [21], robotics [15], learning algorithm [29], and systems biology [16], [20].

Any two solutions of a contractive system converge to each other, which implies a unique exponentially asymptotically stable equilibrium or trajectory. This rules out the existence of multiple (stable or unstable) equilibriums, limit cycles, and other oscillatory behaviors. This motivates researchers to introduce generalizations of contraction theory which allow analyzing non-trivial attractors, for example, partial contraction [28], [7], [24], horizontal contraction [10], and $k$-contraction [30].

Roughly speaking, partial contraction is related to the contractive behavior of an auxiliary system associated with the studied one, and horizontal contraction studies contractive properties along some particular "directions". Despite using different mathematical formulations, these two generalized

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Fig. 1. Illustration of the main results.
notions of contraction are both effective for analyzing stable limit cycles or synchronization of networked systems, i.e., convergence to certain subspaces. Furthermore, Ref. [27] showed that both notions (where partial contraction is referred to as virtual contraction) can be utilized to solve a particular control problem arising in the immersion \& invariance (I\&I) stabilization procedure [4].

Ref. [30] introduced a generalization of contraction theory called $k$-contraction based on the seminal work of Muldowney [17]. A dynamical system is called $k$-contractive if the dynamics contracts $k$-volumes at an exponential rate. For $k=1$, this recovers the standard contraction theory as 1 volume is just length. However, $k$-contraction with $k>1$ can be applied to analyze systems that are not contractive (i.e., not 1-contractive) such as multi-stable systems that are prevalent in mathematical models of real-world systems. In particular, it was shown in [30], [31] that $k$-contraction can also be applied to chaotic systems, which typically cannot be analyzed using partial/horizontal contraction.

However, the three notions of contraction are closely related in certain cases and it is our intention to bring more insights into the distinctions and the relations. Unlike [30], this current work considers dynamical systems whose solutions evolve on a forward invariant and connected, but not necessarily convex manifold. In our main results, we provide conditions describing when partial contraction implies horizontal contraction and when horizontal contraction implies $k$-contraction. By combining these results together, sufficient conditions for partial contraction implying $k$-contraction are also obtained (see Fig. 1). Furthermore, some examples are given to show that these three notions are different in general. These results are useful since $k$-contraction and, in particular, 2 -contraction implies strong results on the attractors and asymptotic behavior of nonlinear time-invariant systems [11], [30]. Therefore, the same conclusions can be drawn for partially or horizontally contractive systems if the aforementioned sufficient conditions hold.

The remainder of this note is organized as follows. The next section briefly reviews the definitions of $k$-contraction, horizontal contraction, and partial contraction. Section III and IV detail the main results. The Andronov-Hopf oscillator is revisited in Section V to validate the proposed results. Section VI gives the conclusions.

Notation. Here we briefly describe the basic notations and some mathematical tools including compound matrices and wedge products are required to define $k$-contraction, see [5], [17], [9] for more details and proofs.

For two integers $i, j$, with $0<i \leq j$, we denote $[i, j]:=$ $\{i, i+1, \ldots, j\}$. For an $n$-dimensional manifold $\Omega \subseteq \mathbb{R}^{n}$, we denote the tangent space of $\Omega$ at $x \in \Omega$ by $T_{x} \Omega$, and the tangent bundle of $\Omega$ by $T \Omega:=\cup_{x \in \Omega}\left(\{x\} \times T_{x} \Omega\right)$.

Compound matrices. Given $A \in \mathbb{R}^{n \times m}$ and $k \in$ $[1, \min \{n, m\}]$, the $k$ th multiplicative compound matrix of $A$, denoted $A^{(k)}$, is the $\binom{n}{k} \times\binom{ m}{k}$ matrix that includes all the minors of order $k$ of $A$ ordered lexicographically. In particular, $A^{(1)}=A$, and for $A \in \mathbb{R}^{n \times n}, A^{(n)}=\operatorname{det}(A)$. The Cauchy-Binet formula (see e.g., [8, Ch. 1]) asserts that for any $A \in \mathbb{R}^{n \times p}, B \in \mathbb{R}^{p \times m}$, and $k=[1, \min \{n, m, p\}]$,

$$
\begin{equation*}
(A B)^{(k)}=A^{(k)} B^{(k)} \tag{1}
\end{equation*}
$$

which justifies the term multiplicative compound.
The $k$ th additive compound matrix of a square matrix $A \in$ $\mathbb{R}^{n \times n}$ is defined by

$$
A^{[k]}:=\left.\frac{\mathrm{d}}{\mathrm{~d} \varepsilon}\left(I_{n}+\varepsilon A\right)^{(k)}\right|_{\varepsilon=0}
$$

where $I_{n}$ denotes the $n \times n$ identity matrix. This implies that

$$
\begin{equation*}
\left(I_{n}+\varepsilon A\right)^{(k)}=I_{r}+\varepsilon A^{[k]}+o(\varepsilon) \tag{2}
\end{equation*}
$$

with $r:=\binom{n}{k}$. The matrix $A^{[k]}$ can be given explicitly in terms of the entries $a_{i j}$ of $A$ as shown in [23], [9]. It follows from (2) and the properties of the multiplicative compound that for any $A, B \in \mathbb{R}^{n \times n}$

$$
\begin{equation*}
(A+B)^{[k]}=A^{[k]}+B^{[k]} \tag{3}
\end{equation*}
$$

which justifies the term additive compound.
Wedge products. The multiplicative compound matrix has an important geometric interpretation in terms of the $k$-volume of a $k$-parallelotope [5]. Pick $k \in[1, n]$, and $k$ vectors $a^{i} \in \mathbb{R}^{n}$, $i=1, \ldots, k$. The wedge product of these vectors, denoted $a^{1} \wedge$ $\cdots \wedge a^{k}$, can be represented using the multiplicative compound as

$$
a^{1} \wedge \cdots \wedge a^{k}=\left[\begin{array}{lll}
a^{1} & \ldots & a^{k} \tag{4}
\end{array}\right]^{(k)}
$$

This provides a representation of wedge product as an $r$ dimensional column vector, where $r:=\binom{n}{k}$. We will use the short-hand notation $\wedge_{i=1}^{k} a^{i}:=a^{1} \wedge \cdots \wedge a^{k}$ throughout the paper. By the property of the multiplicative compound matrices, $\wedge_{i=1}^{k} a^{i}=0$ iff $a^{1}, \ldots, a^{k}$ are linearly dependent [17]. The $k$-parallelotope generated by $a^{1}, \ldots, a^{k}$ (and the zero vertex) is $P\left(a^{1}, \ldots, a^{k}\right):=\left\{\sum_{i=1}^{k} c_{i} a^{i} \mid c_{i} \in[0,1]\right\}$. The $k$-volume of $P\left(a^{1}, \ldots, a^{k}\right)$ is $\left|\wedge_{i=1}^{k} a^{i}\right|_{2}$, where $|\cdot|_{2}$ is the $L_{2}$ norm. In the particular case $k=n$ this reduces to the wellknown formula: volume $\left(P\left(a^{1}, \ldots, a^{n}\right)\right)=\left|\operatorname{det}\left(a^{1}, \ldots, a^{n}\right)\right|$.

## II. Three Generalized Notions of Contraction

In this section, we briefly review the definitions of $k$ contraction, horizontal contraction, and partial contraction.

Consider the nonlinear time-varying (NTV) system

$$
\begin{equation*}
\dot{x}=f(t, x) \tag{5}
\end{equation*}
$$

where $f: \mathbb{R}_{+} \times \mathbb{R}^{n} \rightarrow \mathbb{R}^{n}$ is continuously differentiable w.r.t. its second argument, and let $J(t, x):=\frac{\partial}{\partial x} f(t, x)$ denote the Jacobian of the vector field w.r.t. $x$. We denote by $x\left(t, t_{0}, a\right)$ the solution to (5) at time $t$ emanating from the initial condition $a \in \mathbb{R}^{n}$ at time $t_{0}$, that is, $x\left(t_{0}, t_{0}, a\right)=a$. We assume that the solutions of (5) evolve on a closed and connected $n$-dimensional manifold $\Omega$, and that for any initial condition $a \in \Omega$, a unique solution $x\left(t, t_{0}, a\right)$ exists and satisfies $x\left(t, t_{0}, a\right) \in \Omega$ for all $t \geq t_{0}$. For the sake of simplicity, we assume from here on that the initial time is $t_{0}=0$, and write $x(t, a):=x(t, 0, a)$.

Consider the matrix $\Phi(t, a):=\frac{\partial x(t, a)}{\partial a}$. Note that $\Phi(0, a)=$ $I_{n}$. A straightforward computation yields

$$
\begin{equation*}
\frac{d}{d t} \Phi(t, a)=J(t, x(t, a)) \Phi(t, a) \tag{6}
\end{equation*}
$$

This is the variational system associated with (5) along $x(t, a)$. Let $\delta a \in T_{a} \Omega$ denote an infinitesimal variation to the initial condition $a \in \Omega$. Then $\delta x(t, a):=\Phi(t, a) \delta a$ is the infinitesimal displacement w.r.t. the solution $x(t, a)$ induced by the initial condition $a+\delta a$. Eq. (6) implies that

$$
\begin{equation*}
\delta \dot{x}(t, a):=\frac{\mathrm{d}}{\mathrm{~d} t} \delta x(t, a)=J(t, x(t, a)) \delta x(t, a) \tag{7}
\end{equation*}
$$

## A. $k$-contraction

For $k \in[1, n]$, consider $\Phi^{(k)}(t, a):=(\Phi(t, a))^{(k)}$, i.e., the $k$ th multiplicative compound matrix of $\Phi(t, a)$. Fix $\varepsilon>0$. Eqs. (6) and (1) give

$$
\begin{aligned}
\Phi^{(k)}(t+\varepsilon, a) & =(\Phi(t, a)+\varepsilon J(t, x(t, a)) \Phi(t, a))^{(k)}+o(\varepsilon) \\
& =\left(I_{n}+\varepsilon J(t, x(t, a))\right)^{(k)} \Phi^{(k)}(t, a)+o(\varepsilon)
\end{aligned}
$$

Combining this with (2) and the fact that $\Phi(0, a)=I_{n}$ yields a differential equation for $\Phi^{(k)}(t, a)$ :

$$
\begin{align*}
\frac{d}{d t} \Phi^{(k)}(t, a) & =\lim _{\varepsilon \rightarrow 0^{+}} \frac{\Phi^{(k)}(t+\varepsilon, a)-\Phi^{(k)}(t, a)}{\varepsilon} \\
& =J^{[k]}(t, x(t, a)) \Phi^{(k)}(t, a), \Phi^{(k)}(0, a)=I_{r} \tag{8}
\end{align*}
$$

where $J^{[k]}(t, x(t, a)):=(J(t, x(t, a)))^{[k]}$, and $r:=\binom{n}{k}$. In other words, all the minors of order $k$ of $\Phi(t, a)$, stacked in the matrix $\Phi^{(k)}(t, a)$, satisfy a linear dynamics with the matrix $J^{[k]}(t, x(t, a))$.

Pick $k$ initial conditions $\delta a^{1}, \ldots, \delta a^{k}$ for (7). Define

$$
\begin{equation*}
\delta x^{i}(t, a):=\Phi(t, a) \delta a^{i}, \quad y(t, a):=\wedge_{i=1}^{k} \delta x^{i}(t, a) \tag{9}
\end{equation*}
$$

Note that $y(t, a)=0$ iff $\delta x^{1}(t, a), \ldots, \delta x^{k}(t, a)$ are linearly dependent. By (1) and (4),

$$
\begin{equation*}
y(t, a)=\wedge_{i=1}^{k} \Phi(t, a) \delta a^{i}=\Phi^{(k)}(t, a) y(0, a) \tag{10}
\end{equation*}
$$

and (8) yields

$$
\begin{equation*}
\dot{y}(t, a)=J^{[k]}(t, x(t, a)) y(t, a) \tag{11}
\end{equation*}
$$

This is the $k$ th compound equation of (7) along $x(t, a)$ (see e.g. [12]). This leads to the following definition.

Definition 1 ( $k$-contraction). Fix $k \in[1, n]$ and let $r:=\binom{n}{k}$. The NTV system (5) is called $k$-contractive if the linear timevarying (LTV) system

$$
\begin{equation*}
\dot{y}(t)=J^{[k]}(t, x(t, a)) y(t) \tag{12}
\end{equation*}
$$

is uniformly exponentially stable for any $a \in \Omega$, that is, there exist $c \geq 1, \eta>0$, and a vector norm $|\cdot|$ such that

$$
\begin{equation*}
|y(t)| \leq c \exp (-\eta t)|y(0)|, \text { for all } t \geq 0 \tag{13}
\end{equation*}
$$

Note that the above definition is slightly different from [30, Def. 2] where the value of $c$ is fixed as one. By the geometric interpretation of wedge products, Eq. (13) implies that the $k$-volume of the $k$-parallelotope generated by the vertices $\delta x^{1}(t, a), \ldots, \delta x^{k}(t, a)$ (and the zero vertex) decays to zero at an exponential rate. For $k=1, k$-contraction reduces to standard contraction.

## B. Horizontal contraction

For every $x \in \Omega$, suppose that $T_{x} \Omega$ can be subdivided into a horizontal distribution $\mathcal{H}_{x}$ and a vertical distribution $\mathcal{Q}_{x}$ which are orthogonally complementary to each other. That is, there exist $\ell \in[1, n]$ and $C^{1}$ mappings $h^{i}, q^{i}: \mathbb{R}^{n} \rightarrow \mathbb{R}^{n}$ such that

$$
\begin{align*}
\mathcal{H}_{x} & :=\operatorname{span}\left\{h^{1}(x), \ldots, h^{\ell}(x)\right\} \\
\mathcal{Q}_{x} & :=\operatorname{span}\left\{q^{1}(x), \ldots, q^{n-\ell}(x)\right\} \tag{14}
\end{align*}
$$

Note that if $\ell=n$, then $\mathcal{H}_{x}=T_{x} \Omega$. Define the matrices

$$
\begin{align*}
& H(x):=\left[\begin{array}{lll}
h^{1}(x) & \cdots & h^{\ell}(x)
\end{array}\right] \in \mathbb{R}^{n \times \ell}  \tag{15}\\
& Q(x):=\left[\begin{array}{lll}
q^{1}(x) & \cdots & q^{n-\ell}(x)
\end{array}\right] \in \mathbb{R}^{n \times(n-\ell)}
\end{align*}
$$

Since $\mathcal{H}_{x}$ and $\mathcal{Q}_{x}$ are orthogonal to each other, we have

$$
\begin{equation*}
H^{T}(x) Q(x)=0 \tag{16}
\end{equation*}
$$

For every $\delta x \in T_{x} \Omega$, there exists a set of uniquely defined $\delta x_{h} \in \mathbb{R}^{\ell}$ and $\delta x_{q} \in \mathbb{R}^{n-\ell}$ such that

$$
\begin{equation*}
\delta x=H(x) \delta x_{h}+Q(x) \delta x_{q} . \tag{17}
\end{equation*}
$$

Note that $H(x) \delta x_{h} \in \mathcal{H}_{x}$, and $Q(x) \delta x_{q} \in \mathcal{Q}_{x}$. Combining (17) and (16) gives

$$
\begin{align*}
& H^{T}(x) \delta x=H^{T}(x) H(x) \delta x_{h} \\
& Q^{T}(x) \delta x=Q^{T}(x) Q(x) \delta x_{q}, \text { for all }(x, \delta x) \in T \Omega \tag{18}
\end{align*}
$$

Without loss of generality, we assume throughout that both $H(x)$ and $Q(x)$ are bounded on $x \in \Omega$. Based on the above discussions, horizontal contraction is defined as follows.
Definition 2 (Horizontal contraction). The NTV system (5) is called horizontally contractive w.r.t. $\mathcal{H}_{x}$ if there exist $c \geq$ $1, \eta>0$, and a vector norm $|\cdot|$ such that the solution of (7) for any $a \in \Omega$, i.e., $\delta x(t, a)=H(x(t, a)) \delta x_{h}(t, a)+$ $Q(x(t, a)) \delta x_{q}(t, a)$, satisfies

$$
\begin{equation*}
\left|\delta x_{h}(t, a)\right| \leq c \exp (-\eta t)\left|\delta x_{h}(0, a)\right|, \quad \text { for all } t \geq 0 \tag{19}
\end{equation*}
$$

In [10], horizontal contraction is formalized via a differential Lyapunov framework. Specifically, a sufficient condition is given in terms of a so-called Horizontal Finsler-Lyapunov function.

Definition 3 (Horizontal Finsler-Lyapunov function [10]). Consider a manifold $\Omega$ and the tangent space $T_{x} \Omega=\mathcal{H}_{x} \oplus \mathcal{Q}_{x}$, where $\oplus$ denotes the direct sum of vector spaces. A $C^{1}$ function $V: T \Omega \rightarrow \mathbb{R}_{+}$is called a candidate horizontal Finsler-Lyapunov function for (5) if there exist constants $d_{1}, d_{2}>0, d_{3}>1$, and a function $F: T \Omega \rightarrow \mathbb{R}_{+}$such that

$$
\begin{align*}
& V(x, \delta x)=V\left(x, H(x) \delta x_{h}\right) \\
& d_{1}(F(x, \delta x))^{d_{3}} \leq V(x, \delta x) \leq d_{2}(F(x, \delta x))^{d_{3}}  \tag{20}\\
& \quad \text { for all }(x, \delta x) \in T \Omega
\end{align*}
$$

and $F$ satisfies the following conditions:
(i) $F(x, \delta x)=F\left(x, H(x) \delta x_{h}\right)$ for every $(x, \delta x) \in T_{x} \Omega$;
(ii) $F(x, \delta x)$ is $C^{1}$ for all $x \in \Omega$ and $\delta x \in \mathcal{H}_{x} \backslash\{0\}$;
(iii) $F(x, \delta x) \geq 0$ for all $(x, \delta x) \in T \Omega$ with equality only when $\delta x \in \mathcal{Q}_{x}$;
(iv) $F(x, \lambda \delta x)=|\lambda| F(x, \delta x)$ for all $(x, \delta x) \in T \Omega$ and any $\lambda \in \mathbb{R}$;
(v) $F\left(x, \delta x^{1}+\delta x^{2}\right) \leq F\left(x, \delta x^{1}\right)+F\left(x, \delta x^{2}\right)$ for all $\left(x, \delta x^{1}\right),\left(x, \delta x^{2}\right) \in T \Omega$, with equality only when $\delta x_{h}^{1}=$ $\lambda \delta x_{h}^{2}$ for some $\lambda \in \mathbb{R}$;
(vi) there exist constants $d_{4}, d_{5}>0$ and a vector norm $|\cdot|$ such that $d_{4}\left|\delta x_{h}\right| \leq F(x, \delta x) \leq d_{5}\left|\delta x_{h}\right|$ for all $(x, \delta x) \in$ $T \Omega$.

Proposition 1. [10] Consider the NTV system (5) and the associated variational system (7) with a candidate horizontal Finsler-Lyapunov function $V: T \Omega \rightarrow \mathbb{R}_{+}$. If there exists $\lambda>$ 0 such that

$$
\begin{equation*}
\frac{\partial V(x, \delta x)}{\partial x} f(t, x)+\frac{\partial V(x, \delta x)}{\partial \delta x} J(t, x) \delta x \leq-\lambda V(x, \delta x) \tag{21}
\end{equation*}
$$

for all $t \in \mathbb{R}_{+}$and all $(x, \delta x) \in T \Omega$, then the system (5) is horizontally contractive w.r.t. $\mathcal{H}_{x}$.

By taking the derivative of $V(x(t), \delta x(t))$ along the trajectories of (5) and (7), we have

$$
\begin{equation*}
\dot{V}(x, \delta x) \leq-\lambda V(x, \delta x), \quad \text { for all } t \geq 0 \tag{22}
\end{equation*}
$$

so $V(x(t), \delta x(t)) \leq \exp (-\lambda t) V(x(0), \delta x(0))$. By (20) and Property (vi) in Definition 3, Eq. (19) holds with $c:=$ $\left(d_{2} / d_{1}\right)^{\frac{1}{d_{3}}} d_{5} / d_{4}$ and $\eta:=\lambda / d_{3}$.

If $\mathcal{Q}_{x}=\{0\}$ in Definition 2, i.e., $\mathcal{H}_{x}=T_{x} \Omega$ and $H(x)=$ $I_{n}$, then (19) is the same as (13) in Definition 1 with $k=1$, since $y(t)$ in (13) is a solution of the variational system (7) in this case. Therefore, Definition 2 with $\mathcal{H}_{x}=T_{x} \Omega$ reduces to the definition for standard contraction. In this case, $V(x, \delta x)$ in Definition 3 is called a Finsler-Lyapunov function [10, Def. 2].

## C. Partial contraction

The following definition is a time-varying version of partial contraction as given in [7].

Definition 4 (Partial contraction). Consider the NTV system (5). Assume that there exist functions $p(x): \mathbb{R}^{n} \rightarrow \mathbb{R}^{\ell}$, with $\ell \in[1, n]$, and $g: \mathbb{R}_{+} \times \mathbb{R}^{\ell} \times \mathbb{R}^{n} \rightarrow \mathbb{R}^{n}$ such that

$$
\begin{equation*}
g(t, p(x), x)=f(t, x) \tag{23}
\end{equation*}
$$

System (5) is called partially contractive w.r.t. $p(x)$ if the system

$$
\begin{equation*}
\dot{\xi}=f_{\xi}(t, \xi, x):=\frac{\partial p(x)}{\partial x} g(t, \xi, x) \tag{24}
\end{equation*}
$$

is contractive w.r.t. $\xi$ for all $t \in \mathbb{R}_{+}, x \in \Omega$, and $\xi \in \Omega_{\xi}$. Here, $\Omega_{\xi} \subseteq \mathbb{R}^{\ell}$ denotes the state-space of (24).

Here, the functions $p(\cdot)$ and $g(\cdot)$ are related to a $f a c$ torization of $f(t, x)$ [7]. Partial contraction implies that every solution of (24) converges to $p(x(t))$ exponentially since $\xi(t)=p(x(t))$ is a particular solution of (24). If $\xi(t)=0$ is also a solution of (24), then the contraction property implies that all the trajectories of (5) converge to the manifold $\mathcal{M}:=\left\{x \in \mathbb{R}^{n} \mid p(x)=0\right\}$. For the special case that $p(x)=x$, the system (24) may serve as an observer for (5) [28].

Let $\delta \xi_{0} \in T \Omega_{\xi}$ denote an infinitesimal virtual variation to the initial condition $\xi(0)=\xi_{0} \in \mathbb{R}^{\ell}$ of (24). Define $\delta \xi(t):=\frac{\partial \xi\left(t, \xi_{0}\right)}{\partial \xi_{0}} \delta \xi_{0}$. This yields the following variational system associated with (24):

$$
\begin{equation*}
\delta \dot{\xi}(t)=J_{\xi}(t, \xi, x) \delta \xi(t) \tag{25}
\end{equation*}
$$

where $J_{\xi}(t, \xi, x):=\frac{\partial f_{\xi}}{\partial \xi}(t, \xi, x)$. This variational system will be instrumental in the subsequent analysis.

## III. From partial contraction to horizontal CONTRACTION

The next example shows that partial contraction does not necessarily imply horizontal contraction. Roughly speaking, this is due to that the choice of $p(x)$ in Definition 4 is less conservative than the choice of $\mathcal{H}_{x}$ in Definition 2.
Example 1. Consider the following linear system with a nonlinear perturbation

$$
\begin{equation*}
\dot{x}=A x+b(t, x) \tag{26}
\end{equation*}
$$

where $A \in \mathbb{R}^{n \times n}$ is assumed to be Hurwitz, and $b: \mathbb{R}_{+} \times$ $\mathbb{R}^{n} \rightarrow \mathbb{R}^{n}$. Let $p(x)=x$, then (24) in this case can be constructed as: $\dot{\xi}=A \xi+b(t, x)$, which is contractive w.r.t. $\xi$. That is, the system (26) is partially contractive w.r.t. $p(x)=x$.

If we let $b(t, x)=-A x_{d}(t)+\dot{x}_{d}(t)$, and $x_{d}: \mathbb{R}_{+} \rightarrow \mathbb{R}^{n}$ be a time parameterized closed curve with self intersections (e.g., for $n=2$, set $\left.x_{d}(t)=\left[\begin{array}{ll}\sin \left(\frac{2 \pi t}{3}\right) & 1.5 \cos \left(\frac{\pi t}{3}\right)\end{array}\right]^{T}\right)$. Note that $x(t)=x_{d}(t)$ is a system solution in this case. The existence of self intersections in $x_{d}$ implies that this system is not contractive in any specified direction, i.e., it is not horizontally contractive. As another example, let $A=-c I_{3}, c \approx 0.208186$, and $b(t, x)=\left[\begin{array}{lll}\sin \left(x_{2}\right) & \sin \left(x_{3}\right) & \sin \left(x_{1}\right)\end{array}\right]$, then (26) reduces to the chaotic system introduced by René Thomas (see e.g., [31, Sec. V]), which admits a strange attractor with no explicit analytical form. This implies that a horizontal distribution $\mathcal{H}_{x}$ is difficult to construct in this case and possibly does not exist.

The next result specifies a sufficient condition such that partial contraction implies horizontal contraction.

Theorem 1. Suppose that there exists $\lambda>0$, and a FinslerLyapunov function $V_{\xi}: T \Omega_{\xi} \rightarrow \mathbb{R}_{+}$for the system (24) such that
$\frac{\partial V_{\xi}(\xi, \delta \xi)}{\partial \xi} f_{\xi}(t, \xi, x)+\frac{\partial V_{\xi}(\xi, \delta \xi)}{\partial \delta \xi} J_{\xi}(t, \xi, x) \delta \xi \leq-\lambda V_{\xi}(\xi, \delta \xi)$,
for all $t \in \mathbb{R}_{+}, x \in \Omega$, and $(\xi, \delta \xi) \in T \Omega_{\xi}$. That is, the system (5) is partially contractive w.r.t. $p(x)$. Consider a distribution $\mathcal{H}_{x}$ as in (14) and the associated matrix $H(x) \in$ $\mathbb{R}^{n \times \ell}$ in (15). Assume that there exists a vector norm $|\cdot|$, and $d_{6}, d_{7}>0$ such that

$$
\begin{equation*}
d_{6}|y| \leq\left|H^{T}(x) H(x) y\right| \leq d_{7}|y|, \text { for all } x \in \Omega, y \in \mathbb{R}^{\ell} \tag{28}
\end{equation*}
$$

Furthermore, for all $t \in \mathbb{R}_{+}$and all $x \in \Omega$, assume that

$$
\begin{equation*}
H_{f}^{T}(x)+H^{T}(x) J(t, x)=J_{\xi}(t, p(x), x) H^{T}(x) \tag{29}
\end{equation*}
$$

where

$$
H_{f}(x):=\left[\begin{array}{lll}
\frac{\partial h_{1}(x)}{\partial x} f(t, x) & \cdots & \frac{\partial h_{\ell}(x)}{\partial x} f(t, x)
\end{array}\right] \in \mathbb{R}^{n \times \ell}
$$

is the the directional derivative of $H(x)$ along $f(t, x)$. Then, the system (5) is also horizontally contractive w.r.t. $\mathcal{H}_{x}$.

Proof: Differentiating $H^{T}(x) \delta x(t)$ along the trajectories of (5) and (7) gives

$$
\begin{align*}
\frac{\mathrm{d}}{\mathrm{~d} t}\left(H^{T}(x) \delta x\right) & =\left(H_{f}^{T}(x)+H^{T}(x) J(t, x)\right) \delta x  \tag{30}\\
& =J_{\xi}(t, p(x), x) H^{T}(x) \delta x
\end{align*}
$$

where (29) is used. This implies that $\delta \xi(t):=H^{T}(x(t)) \delta x(t)$ is a trajectory of (25). Recall that $\xi(t)=p(x(t))$ is a solution of (24). Let

$$
\begin{equation*}
V(x, \delta x):=V_{\xi}\left(p(x), H^{T}(x) \delta x(t)\right) \tag{31}
\end{equation*}
$$

By (18), $V(x, \delta x)=V\left(x, H(x) \delta x_{h}\right)$. Eq. (31) implies that the derivative of $V(x, \delta x)$ is equal to the derivative of $V_{\xi}(\xi, \delta \xi)$ along the trajectories $\xi(t)=p(x)$ and $\delta \xi(t)=H^{T}(x) \delta x(t)$, i.e.,

$$
\frac{\mathrm{d}}{\mathrm{~d} t} V(x(t), \delta x(t))=\left.\frac{\mathrm{d}}{\mathrm{~d} t} V_{\xi}(\xi(t), \delta \xi(t))\right|_{\xi=p(x), \delta \xi=H^{T}(x) \delta x}
$$

Specifically,

$$
\begin{align*}
& \frac{\partial V(x, \delta x)}{\partial x} f(t, x)+\frac{\partial V(x, \delta x)}{\partial \delta x} J(t, x) \\
= & \left(\frac{\partial V_{\xi}(\xi, \delta \xi)}{\partial \xi} f_{\xi}(t, \xi, x)\right. \\
& \left.+\frac{\partial V_{\xi}(\xi, \delta \xi)}{\partial \delta \xi} J_{\xi}(t, \xi, x) \delta \xi\right)\left.\right|_{\xi(t)=p(x), \delta \xi(t)=H^{T}(x) \delta x(t)} \\
\leq & -\left.\lambda V_{\xi}(\xi, \delta \xi)\right|_{\xi(t)=p(x), \delta \xi(t)=H^{T}(x) \delta x(t)} \\
= & -\lambda V(x, \delta x) \tag{32}
\end{align*}
$$

In order to further show that $V(x(t), \delta x(t))$ is indeed a horizontal Finsler-Lyapunov function for (5) w.r.t. $\mathcal{H}_{x}$, we need to find an associated function $F(x, \delta x)$ satisfying Definition 3.

Since $V_{\xi}(\xi, \delta \xi)$ is Finsler-Lyapunov function of the sys-
tem (24), there exists an associated function $F_{\xi}: T \Omega_{\xi} \rightarrow \mathbb{R}_{+}$ that satisfies Definition 3 with a zero vertical distribution. Let

$$
F(x, \delta x):=F_{\xi}\left(p(x), H^{T}(x) \delta x\right)
$$

Hence,

$$
d_{4}\left|H^{T}(x) \delta x\right| \leq F(x, \delta x) \leq d_{5}\left|H^{T}(x) \delta x\right|,
$$

for some $d_{4}, d_{5}>0$. By (18), this yields

$$
d_{4}\left|H^{T}(x) H(x) \delta x_{h}\right| \leq F(x, \delta x) \leq d_{5}\left|H^{T}(x) H(x) \delta x\right|
$$

From (28), we have

$$
d_{4} d_{6}\left|\delta x_{h}\right| \leq F(x, \delta x) \leq d_{5} d_{7}\left|\delta x_{h}\right| .
$$

That is, $F(x, \delta x)$ satisfies property (vi) in Definition 3. Furthermore, note that $V(x, \delta x)$ and $F(x, \delta x)$ satisfy condition (20) and properties (i)-(v) in Definition 3 by directly inheriting those from $V_{\xi}(\xi, \delta \xi)$ and $F_{\xi}(\xi, \delta \xi)$. Hence, $V(x, \delta x)$ is indeed a horizontal Finsler-Lyapunov function for (5). By Prop. 1, Eq. (32) ensures that the system (5) is horizontally contractive w.r.t. $\mathcal{H}_{x}$.

The next example shows that condition (29) generally does not hold even when $p(x)$ is linear. As suggested by the next example and Section V, the choice of $H(x)$ in Theorem 1 is closely related to $p(x)$. Specifically, if partial contraction ensures that all the solutions of (5) converge to the manifold $\mathcal{M}:=\left\{x \in \mathbb{R}^{n} \mid p(x)=0\right\}$, then it is reasonable to assume that the system is contractive along the directions orthogonal to $\mathcal{M}$. That is, a candidate for $H(x)$ can be selected as $H(x)=$ $\frac{\partial^{T} p(x)}{\partial x}$.

Example 2 (Convergence to flow-invariant subspaces [19]). Let $H \in \mathbb{R}^{n \times \ell}$ and $Q \in \mathbb{R}^{n \times(n-\ell)}$, such that

$$
\begin{align*}
H^{T} H & =I_{\ell}, \quad Q^{T} Q=I_{n-\ell} \\
H^{T} Q & =0, \quad H H^{T}+Q Q^{T}=I_{n} \tag{33}
\end{align*}
$$

Note that the above conditions hold if the column vectors of $H$ and $Q$, i.e., $h^{1}, \ldots, h^{\ell}, q^{1}, \ldots, q^{n-\ell}$, are orthonormal, which means that they all have Euclidean norm one and are mutually orthogonal. Let $\mathcal{H}$ and $\mathcal{Q}$ denote the column subspaces of $H$ and $Q$, respectively. Assume that the NTV system (5) satisfies

$$
\begin{equation*}
f(t, \mathcal{Q}) \subseteq \mathcal{Q} \tag{34}
\end{equation*}
$$

That is, $\mathcal{Q}$ is flow-invariant. In this case, Eq. (23) naturally holds with $p(x)=H^{T} x$ and

$$
\begin{equation*}
g(t, p(x), x):=f\left(t, H p(x)+Q Q^{T} x\right) \tag{35}
\end{equation*}
$$

Then, the system (24) can be defined as

$$
\begin{equation*}
\dot{\xi}=f_{\xi}(t, \xi, x):=H^{T} g\left(t, H \xi+Q Q^{T} x\right) \tag{36}
\end{equation*}
$$

Note that $\xi(t)=H^{T} x(t)$ and $\xi(t)=0$ are two particular solutions to (36). Therefore, partial contraction w.r.t. $H^{T} x$ ensures that all the trajectories of (5) converge to $\mathcal{Q}$.

Then, we study if the above conditions also imply horizontal contraction. Consider the case when $f(t, x)=0$ for all $x \in \mathcal{Q}$, which ensures that (34) holds. In this case, the NTV system (5) can only be horizontally contractive w.r.t. $\mathcal{H}$. Without loss of generality, we can choose $H(x)$ in Theorem 1 as $H$ here.

Note that $J_{\xi}\left(t, H^{T} x, x\right)=H^{T} J(t, x) H$. Then, condition (29) in this case boils down to

$$
\begin{equation*}
H^{T} J(t, x)=H^{T} J(t, x) H H^{T} \tag{37}
\end{equation*}
$$

which does not hold in general. According to Theorem 1, if the system (5) is partially contractive w.r.t. $H^{T} x$, then (37) ensures that it is also horizontally contractive w.r.t. $\mathcal{H}$. From (33), a sufficient condition to ensure (37) is

$$
\begin{equation*}
H^{T} J(t, x) Q=0 \tag{38}
\end{equation*}
$$

Consider a linear system, i.e., $J(t, x)=A \in \mathbb{R}^{n \times n}$. Then, Eq. (38) holds naturally since (34) implies $A \mathcal{Q} \subseteq \mathcal{Q}$.

In general, for a horizontally contractive system (5), the existence of $p(x)$ as in Definition 4, which depends on the integrability of $H(x)$, is not ensured. That is, horizontal contraction does not necessarily imply partial contraction.

## IV. FROM HORIZONTAL CONTRACTION TO $k$-CONTRACTION

To facilitate the subsequent result, we first prove a useful property of wedge products. Recall that a vector norm $|\cdot|$ : $\mathbb{R}^{n} \rightarrow \mathbb{R}_{+}$is called monotonic if for any $x, y \in \mathbb{R}^{n}$ with $\left|x_{i}\right| \leq$ $\left|y_{i}\right|$ for all $i$, we have $|x| \leq|y|$ [6]. For example, all the $L_{p}$ norms are monotonic.

Lemma 1. Pick $k \in[2, n]$. Consider a set of $k$ timevarying vectors $a^{1}(t), \cdots, a^{k}(t) \in \mathbb{R}^{n}$. Assume that there exist constants $\ell \in[1, k-1], \gamma_{1} \geq 1, \gamma_{2} \geq 1, \beta>0$, and a monotonic vector norm $|\cdot|$ such that

$$
\begin{align*}
\left|a^{j}(t)\right| & \leq \gamma_{1} \exp (-\beta t)\left|a^{j}(0)\right|, & & j=1, \ldots, \ell  \tag{39}\\
\left|a^{\ell+i}(t)\right| & \leq \gamma_{2}\left|a^{\ell+i}(0)\right|, & & i=1, \ldots, k-\ell \tag{40}
\end{align*}
$$

for all $t \in \mathbb{R}_{+}$. Then, $\left|\wedge_{j=1}^{k} a^{j}(t)\right|$ decays to zero exponentially. Furthermore, for $\left|\wedge_{j=1}^{k} a^{j}(0)\right| \neq 0$, there exists $\bar{\gamma}, \bar{\beta}>0$ such that

$$
\begin{equation*}
\left|\wedge_{j=1}^{k} a^{j}(t)\right| \leq \bar{\gamma} \exp (-\bar{\beta} t)\left|\wedge_{j=1}^{k} a^{j}(0)\right|, \text { for all } t \in \mathbb{R}_{+} . \tag{41}
\end{equation*}
$$

Proof: Recall that $\left|\wedge_{j=1}^{k} a^{j}(t)\right|$ is the $k$-volume of the $k$ parallelotope generated by $a^{1}(t), \cdots, a^{k}(t)$. Intuitively speaking, (39) implies that at least one edge of this $k$-parallelotope shrinks exponentially, and the other edges are uniformly bounded. Therefore, its $k$-volume also shrinks exponentially. Here, we only provide a detailed proof for $\ell=k-1$. The proof for other cases is based on similar arguments.

By the property of the wedge products, i.e., multiplicative compound matrices, we have

$$
\wedge_{j=1}^{k} a^{j}=\sum_{i=1}^{n} a_{i}^{k}\left[\begin{array}{llll}
a^{1} & \ldots & a^{k-1} & e^{i} \tag{42}
\end{array}\right]^{(k)} .
$$

where $a_{i}^{k}$ denotes the $i$ th entry of $a^{k}$, and $e^{i}$ the $i$ th canonical vector in $\mathbb{R}^{n}$. Let $z^{i}(t):=\left[\begin{array}{llll}a^{1} & \ldots & a^{k-1} & e^{i}\end{array}\right]^{(k)}$. Note that $z^{i}(t) \in \mathbb{R}^{r}, r:=\binom{n}{k}$, whose every entry is a minor of order $k$ of $\left[\begin{array}{llll}a^{1} & \ldots & a^{k-1} & e^{i}\end{array}\right]$. By the Leibniz formula for determinants, $z^{i}(t)$ has at least $\binom{n-1}{k}$ zero entries, and every nonzero entry is an entry of the vector $\wedge_{j=1}^{k-1} a^{j}$ multiplied by
either plus one or minus one. Hence, for any monotonic vector norm $|\cdot|: \mathbb{R}^{n} \rightarrow \mathbb{R}_{+}$, we have $\left|z^{i}(t)\right| \leq\left|\wedge_{j=1}^{k-1} a^{j}(t)\right|$ and thus

$$
\begin{aligned}
\left|\wedge_{j=1}^{k} a^{j}(t)\right| & \leq \sum_{i=1}^{n}\left(\left|a_{i}^{k}\right|\left|\wedge_{j=1}^{k-1} a^{j}(t)\right|\right) \\
& \leq n \gamma_{2}\left|a^{k}(0)\right|\left|\wedge_{j=1}^{k-1} a^{j}(t)\right|
\end{aligned}
$$

where we used the fact $\left|a_{i}^{k}(t)\right| \leq\left|a^{k}(t)\right| \leq \gamma_{2}\left|a^{k}(0)\right|$. Then, by repetitively using the recursive formula (42), we have

$$
\left|\wedge_{j=1}^{k} a^{j}(t)\right| \leq\left(n \gamma_{1}\right)^{k-1} \gamma_{2} \exp (-(k-1) \beta t) \prod_{j=1}^{k}\left|a^{j}(0)\right|
$$

That is, $\left|\wedge_{j=1}^{k} a^{j}(t)\right|$ converges to zero exponentially. For the case $\left|\wedge_{j=1}^{k} a^{j}(0)\right| \neq 0$, note that there exists a large enough constant $\gamma_{3}>0$ such that $\prod_{j=1}^{k}\left|a^{j}(0)\right| \leq \gamma_{3}\left|\wedge_{j=1}^{k} a^{j}(0)\right|$. Hence, $\left|\wedge_{j=1}^{k} a^{j}(t)\right| \leq \bar{\gamma} \exp (-\bar{\beta} t)\left|\wedge_{j=1}^{k} a^{j}(0)\right|$, where $\bar{\gamma}:=$ $\left(n \gamma_{1}\right)^{k-1} \gamma_{2} \gamma_{3}$ and $\bar{\beta}:=(k-1) \beta$.

In the subsequent analysis, we assume that the vector norms under concern are monotonic. The next result specifies a sufficient condition such that horizontal contraction implies $k$-contraction. Intuitively speaking, it claims that if a $n$ dimensional horizontally contractive system contracts length in $n-k+1$ directions exponentially, and does not expand length in the remaining $k-1$ directions, then it contracts $k$ volume exponentially.

Theorem 2. Suppose that the system (5) is horizontally contractive w.r.t. $\mathcal{H}_{x}$ defined in (14), where we can write $\ell=n-k+1$ for some $k=[1, n]$. For the matrices $H(x) \in \mathbb{R}^{n \times(n-k+1)}$ and $Q(x) \in \mathbb{R}^{n \times(k-1)}$ given in (18), define

$$
M(x):=\left[\begin{array}{c}
H^{T}(x)  \tag{43}\\
Q^{T}(x)
\end{array}\right] \in \mathbb{R}^{n \times n}
$$

Assume that for some vector norm $|\cdot|$, there exist $c_{i}>0$, $i=1, \ldots, 6$, such that for all $x \in \Omega$

$$
\begin{align*}
& c_{1}|y| \leq\left|H^{T}(x) H(x) y\right| \leq c_{2}|y|, y \in \mathbb{R}^{n-k+1} \\
& c_{3}|y| \leq|M(x) y| \leq c_{4}|y|, y \in \mathbb{R}^{n}  \tag{44}\\
& c_{5}|y| \leq\left|M^{(k)}(x) y\right| \leq c_{6}|y|, y \in \mathbb{R}^{\binom{n}{k}}
\end{align*}
$$

Furthermore, for any two trajectories of the system (5), denoted $x^{1}(t)$ and $x^{2}(t)$, there exists a constant $\gamma_{1}>1$ such that

$$
\begin{equation*}
\left|x^{1}(t)-x^{2}(t)\right| \leq \gamma_{1}\left|x^{1}(0)-x^{2}(0)\right|, \quad \text { for all } t \geq 0 \tag{45}
\end{equation*}
$$

Then, the system (5) is also $k$-contractive on $\Omega$.
Proof: First, we show that the variational system (7) is uniformly bounded under the condition (45). Pick $\varepsilon>0$, and $x_{0}, z_{0} \in \mathbb{R}^{n}$. Define $z(t, \varepsilon):=\frac{x\left(t, x_{0}+\varepsilon z_{0}\right)-x\left(t, x_{0}\right)}{\varepsilon}$. By (45), we have

$$
\begin{equation*}
|z(t, \varepsilon)| \leq \gamma_{1}\left|z_{0}\right| \tag{46}
\end{equation*}
$$

Since the solutions of (5) are continuously dependent on the initial conditions, the limit $z(t):=\lim _{\varepsilon \rightarrow 0} z(t, \varepsilon)$ exists and $z(t)=\frac{\partial x\left(t, x_{0}\right)}{\partial x_{0}} z_{0}$. This implies that $z(t)$ is a solution of the variational system (7) with the initial condition $z(0)=z_{0}$.

Thus, (46) reduces to

$$
\begin{equation*}
|z(t)| \leq \gamma_{1}\left|z_{0}\right| \tag{47}
\end{equation*}
$$

That is, the variational system (7) is uniformly bounded.
Since the system (5) is horizontally contractive, Definition 2 implies that there exist $\gamma_{2} \geq 1$ and $\beta>0$ such that

$$
\begin{equation*}
\left|\delta x_{h}(t)\right| \leq \gamma_{2} \exp (-\beta t)\left|\delta x_{h}(0)\right|, \quad \text { for all } t \geq 0 \tag{48}
\end{equation*}
$$

Recall that any $\delta x \in T_{x} \Omega$ can be rewritten as $\delta x=$ $H(x) \delta x_{h}+Q(x) \delta x_{q}$, where $\delta x_{h} \in \mathbb{R}^{n-k+1}$, and $\delta x_{q} \in \mathbb{R}^{k-1}$. Hence,

$$
M(x) \delta x=\left[\begin{array}{c}
H^{T}(x) H(x) \delta x_{h}  \tag{49}\\
Q^{T}(x) Q(x) \delta x_{q}
\end{array}\right],
$$

where we use the fact that $H^{T}(x) Q(x)=0$. From (44) and (48), we have

$$
\begin{align*}
\left|H^{T}(x(t)) H(x(t)) \delta x_{h}(t)\right| & \leq c_{2}\left|\delta x_{h}(t)\right|  \tag{50}\\
& \leq c_{2} \gamma_{2} \exp \left(-\beta_{1} t\right)\left|\delta x_{h}(0)\right|
\end{align*}
$$

for all $t \geq 0$. That is, the first $(n-k+1)$ entries of $M(x) \delta x$ converge to zero exponentially. By virtue of monotonic vector norms,

$$
\begin{align*}
\left|Q^{T}(x(t)) Q(x(t)) \delta x_{q}(t)\right| & \leq|M(x(t)) \delta x(t)| \leq c_{4}|\delta x(t)| \\
& \leq c_{4} \gamma_{1}|\delta x(0)|, \text { for all } t \geq 0 \tag{51}
\end{align*}
$$

where (44) and (47) are used. That is, the last $(k-1)$ entries of $M(x) \delta x$ are uniformly bounded.

Then, as defined in (9), consider $k$ trajectories of the variation system (7) specified by $x(t, a)$ with $a \in \Omega$, i.e., $\delta^{1} x(t, a), \ldots, \delta^{k} x(t, a)$, such that $\wedge_{i=1}^{k} \delta^{i} a \neq 0$. Let
$A(t):=\left[\begin{array}{lll}M(x(t, a)) \delta x^{1}(t, a) & \cdots & M(x(t, a)) \delta x^{k}(t, a)\end{array}\right]^{T}$.
Note that $A \in \mathbb{R}^{k \times n}$. Using (1) and (4), we have

$$
\begin{align*}
\left(A^{T}(t)\right)^{(k)} & =\wedge_{i=1}^{k} M(x(t, a)) \delta x^{i}(t, a)  \tag{52}\\
& =M^{(k)}(x(t, a)) y(t, a),
\end{align*}
$$

where $y(t, a):=\wedge_{i=1}^{k} \delta x^{i}(t, a)$ as in (9). Let $a^{i}(t)$ denote the $i$ th column of $A(t)$. From (50) and (51), $\left|a^{i}(t)\right|$ with $i=$ $1, \ldots, n-k+1$, converges to zero exponentially, and $\left|a^{i}(t)\right|$ with $i=n-k+2, \ldots, n$, are uniformly bounded.

Let $Q_{k, n}$ denote the set of increasing sequences of $k$ numbers from $[1, n]$ ordered lexicographically. With a slight abuse of notation, we treat such ordered sequences as sets. The cardinality of $Q_{k, n}$ is $r:=\binom{n}{k}$. The $j$ th element of $Q_{k, n}$ is denoted $\kappa_{j}$. For example, $Q_{2,3}=\left\{\kappa_{1}, \kappa_{2}, \kappa_{3}\right\}$, with $\kappa_{1}=$ $\{1,2\}, \kappa_{2}=\{1,3\}, \kappa_{3}=\{2,3\}$. Then, $\wedge_{i \in \kappa_{j}} a^{i}(t)$ is the $j$ th entry of $M^{(k)}(x(t, a)) y(t, a)$. Note that any $k$ vectors $a^{i}(t), i \in \kappa_{j}$, satisfy condition (39) in Lemma 1. Hence, it is ensured that $\wedge_{i \in \kappa_{j}} a^{i}(t), j=1, \cdots, r$, i.e., all the entries of $M^{(k)}(x(t, a)) y(t, a)$, converge to zero exponentially. Therefore, there exist $\gamma_{3} \geq 1$ and $\eta>0$ such that

$$
\begin{equation*}
\left|M^{(k)}(x(t, a)) y(t, a)\right| \leq \gamma_{3} \exp (-\eta t)\left|M^{(k)}(x(0, a)) y(0, a)\right|, \tag{53}
\end{equation*}
$$

for all $t \geq 0$. Then from (44), we have

$$
\begin{equation*}
|y(t, a)| \leq \frac{c_{6} \gamma_{3}}{c_{5}} \exp (-\eta t)|y(0, a)| \tag{54}
\end{equation*}
$$



Fig. 2. Trajectories of the Duffing oscillator for two different initial conditions (black circle).
for all $t \geq 0$. This implies that the system (5) is $k$-contractive by Definition 1 .

Remark 1. Condition (45) actually asserts uniform incremental stability, that is, every solution of (5) is uniformly stable. By Coppel's inequality (see e.g., [26]), condition (45) is ensured if there exists a matrix measure $\mu: \mathbb{R}^{n \times n} \rightarrow \mathbb{R}$ and a constant $c>0$ such that $\int_{0}^{t} \mu(J(s, x(s, a))) \mathrm{d} s \leq$ $c$, for all $a \in \Omega, t \geq 0$. As shown in the proof of Theorem 2 , condition (45) can actually be relaxed since it is only required that the exponential convergence rate of $H^{T}(x) H(x) \delta x_{h}$ is larger than the expansion rate of $Q^{T}(x) Q(x) \delta x_{q}$ to conclude that $y(t, a)$ converges to zero exponentially. For example, consider the linear system

$$
\begin{aligned}
& \dot{x}_{1}=c_{1} x_{1}+x_{2} \\
& \dot{x}_{2}=-c_{2} x_{2}
\end{aligned}
$$

where $c_{2}>c_{1}>0$. This system is 2 -contractive, and horizontally contractive w.r.t. $\mathcal{H}_{x}:=\operatorname{span}\left\{\left[\begin{array}{ll}0 & 1\end{array}\right]^{T}\right\}$. However, it is not uniformly stable since $x_{1}(t)$ will go to infinity. Note also that this system is still horizontally contractive but not 2 -contractive with $c_{1}>c_{2}>0$.

Since $k$-contraction does not require the existence of the distributions $\mathcal{H}_{x}$ and $\mathcal{Q}_{x}$. Therefore, $k$-contraction does not ensure horizontal contraction in general. This is shown by the next example.

Example 3 (The forced Duffing oscillator [25]). Consider the NTV system

$$
\begin{align*}
& \dot{x}_{1}=x_{2} \\
& \dot{x}_{2}=\theta_{1} x_{1}-\theta_{2} x_{1}^{3}-\theta_{3} x_{2}+\theta_{4} \cos \left(\theta_{5} t\right) \tag{55}
\end{align*}
$$

where $\theta_{i}>0, i=1, \ldots, 5$. For $\theta_{1}=\theta_{2}=1, \theta_{3}=0.3, \theta_{4}=$ 0.37 and $\theta_{5}=1.2$, this system has an attractor with self intersections as shown in Fig. 2. Therefore, this system is not contractive in any specified direction, that is, it is not horizontally contractive. Indeed, it can have more complicated attractors with other different parameters. However, the Jacobian $J(t, x)$ of (55), satisfies $J^{[2]}(t, x)=-\theta_{3}<0$. Therefore, it is 2-contractive according to Definition 1 and [30, Thm. 4].

## V. An Example: the Andronov-Hopf oscillator

In this section, the Andronov-Hopf oscillator is revisited to illustrate our results. It is shown that this system is partially contractive, horizontally contractive, and 2-contractive.

The dynamics of the Andronov-Hopf oscillator is

$$
\begin{align*}
& \dot{x}_{1}=-x_{2}-x_{1}\left(x_{1}^{2}+x_{2}^{2}-1\right) \\
& \dot{x}_{2}=x_{1}-x_{2}\left(x_{1}^{2}+x_{2}^{2}-1\right) \tag{56}
\end{align*}
$$

Note that $x=0$ is an unstable equilibrium, and the unit circle is a stable limit cycle. The associated variational system is

$$
\begin{equation*}
\delta \dot{x}(t)=J(x(t)) \delta x(t) \tag{57}
\end{equation*}
$$

where $J(x)=\left[\begin{array}{cc}1-3 x_{1}^{2}-x_{2}^{2} & -2 x_{1} x_{2}-1 \\ -2 x_{1} x_{2}+1 & 1-x_{1}^{2}-3 x_{2}^{2}\end{array}\right]$.
Proposition 2. Fix $0<\gamma_{1}<1<\gamma_{2}$. Consider the manifold $\Omega:=\left\{x \in \mathbb{R}^{2} \mid \gamma_{1} \leq x_{1}^{2}+x_{2}^{2} \leq \gamma_{2}\right\}$. The system (56) is partially contractive, horizontally contractive, and 2 -contractive on $\Omega$.

Proof: Note that $\Omega$ here is forward invariant and connected, but not convex.
(i) Partial contraction: Let $p(x):=x_{1}^{2}+x_{2}^{2}-1$, and $g(\xi, x):=\left[\begin{array}{ll}-x_{2}-x_{1} \xi & x_{1}-x_{2} \xi\end{array}\right]^{T}$. Then $g(p(x), x)=$ $f(x)$. That is, condition (23) in Definition 4 is satisfied. In this case, the system (24) reduces to

$$
\begin{equation*}
\dot{\xi}(t)=-2\left(x_{1}^{2}(t)+x_{2}^{2}(t)\right) \xi(t) \tag{58}
\end{equation*}
$$

A Finsler-Lypapunov function for (58) can be selected as $V_{\xi}(\xi, \delta \xi)=\delta \xi^{2}$, and we have $\dot{V}_{\xi}(\xi, \delta \xi) \leq-4 \gamma_{1} V_{\xi}(\xi, \delta \xi)$. Therefore, the system (56) is partially contractive w.r.t. $p(x)=$ $x_{1}^{2}+x_{2}^{2}-1$. Note that both $p(x)$ and the origin are the solutions to (58). Therefore, for any $a \in \Omega$,

$$
\begin{equation*}
|p(x(t, a))| \leq \exp \left(-2 \gamma_{1} t\right)|p(a)|, \text { for all } t \in \mathbb{R}_{+} \tag{59}
\end{equation*}
$$

(ii) From partial contraction to horizontal contraction: Let $H(x)=\frac{1}{x_{1}^{2}+x_{2}^{2}}\left[\begin{array}{l}x_{1} \\ x_{2}\end{array}\right]$. Note that condition (28) in Theorem 1 holds since $\gamma_{2}^{-1} \leq H^{T}(x) H(x)=\left(x_{1}^{2}+x_{2}^{2}\right)^{-1}<$ $\gamma_{1}^{-1}$, for all $x \in \Omega$. Based on a straightforward calculation,

$$
\frac{\mathrm{d}}{\mathrm{~d} t}\left(H^{T}(x(t)) \delta x(t)\right)=-2\left(x_{1}^{2}(t)+x_{2}^{2}(t)\right) H^{T}(x(t)) \delta x(t)
$$

which implies that condition (29) in Theorem 1 also holds. Therefore, Theorem 1 implies that this system is also horizontally contractive w.r.t. $\mathcal{H}_{x}:=\operatorname{span}\left\{\frac{1}{x_{1}^{2}+x_{2}^{2}}\left[\begin{array}{l}x_{1} \\ x_{2}\end{array}\right]\right\}$. Furthermore, as shown in the proof of Theorem 1, a horizontal Finsler-Lyapunov function can be constructed as:

$$
\begin{aligned}
V(x, \delta x) & =V_{\xi}\left(p(x), H^{T}(x) \delta x\right)=(\delta x)^{T} H(x) H^{T}(x) \delta x \\
& =\left(\frac{x_{1} \delta x_{1}+x_{2} \delta x_{2}}{x_{1}^{2}+x_{2}^{2}}\right)^{2}
\end{aligned}
$$

Indeed, $\dot{V}(x, \delta x)=-4\left(x_{1}^{2}+x_{2}^{2}\right) V(x, \delta x) \leq-4 \gamma_{1} V(x, \delta x)$ for all $x \in \Omega$.
(iii) From horizontal contraction to 2-contraction: The matrix $J_{\text {sym }}(x):=\left(J(x)+J^{T}(x)\right) / 2$ has eigenvalues: $\lambda_{1}(x)=$ $1-x_{1}^{2}-x_{2}^{2}, \lambda_{2}(x)=1-3 x_{1}^{2}-3 x_{2}^{2}$. Hence, $\mu_{2}(J(x))=$
$1-x_{1}^{2}-x_{2}^{2}$, where $\mu_{2}(\cdot)$ is the matrix measure associated with the $L_{2}$ norm. By (59), we have $\int_{0}^{t} \mu_{2}(J(x(s, a))) \mathrm{d} s \leq$ $\frac{\max \left\{\gamma_{2}-1,1-\gamma_{1}\right\}}{2 \gamma_{1}}$, for all $a \in \Omega$ and $t \geq 0$. This implies that condition (45) holds as shown in Remark 1. Let $H(x):=$ $\frac{1}{x_{1}^{2}+x_{2}^{2}}\left[\begin{array}{ll}x_{1} & x_{2}\end{array}\right]^{T}$, and $Q(x):=\frac{1}{x_{1}^{2}+x_{2}^{2}}\left[\begin{array}{ll}-x_{2} & x_{1}\end{array}\right]^{T}$. Condition (44) holds for all $x \in \Omega$ since $\Omega$ is compact. Therefore, Theorem 2 ensures that the system (56) is 2 -contractive.

In order to show this explicitly, for the 2 nd compound equation of (57), i.e., $\dot{y}(t)=J^{[2]}(x(t)) y(t)$, we consider a change of coordinate $y_{m}=M^{(2)}(x) y=\frac{y}{x_{1}^{2}+x_{2}^{2}}$ with $M(x):=\left[\begin{array}{ll}H(x) & Q(x)\end{array}\right]^{T}$. Then,

$$
\begin{aligned}
\dot{y}_{m}(t)= & M^{(2)}(x) J^{[2]}(x)\left(M^{(2)}(x)\right)^{-1} y_{m}(t)+ \\
& M_{f}^{(2)}(x)\left(M^{(2)}(x)\right)^{-1} y_{m}(t) \\
= & -2\left(x_{1}^{2}+x_{2}^{2}\right) y_{m}(t),
\end{aligned}
$$

where $M_{f}^{(2)}(x)$ denotes the directional derivative of $M^{(2)}(x)$ along vector field of (56). This implies that $|y(t)|$ decays to zero exponentially, that is, the system (56) is 2 -contractive on $\Omega$ according to Definition 1 .

## VI. Conclusions

This note shows that partial contraction, horizontal contraction, and $k$-contraction are not equivalent in general. Some sufficient conditions are specified such that $k$-contraction is achieved from partial and horizontal contraction. Since it is known that partial and horizontal contraction can be used to solve synchronization problems of networked systems, a related research direction is to study how $k$-contraction can simplify the analysis of synchronization problems. Another important question is if the specified conditions can become necessary for certain systems.
As shown by [17] and a recent related work [3], $k$ contraction with $k=2$ is very effective to study nonlinear time-invariant systems, in particular, to rule out oscillatory behaviors, and then almost all trajectories converge to an equilibrium (not necessary unique). Therefore, for some partially or horizontally contractive nonlinear time-invariant systems, if the derived conditions hold and lead to 2 -contraction, then we can draw strong conclusions regarding their global asymptotic stability. As a straightforward application, this may simplify the I\&I stabilization method in the sense that the target system there is not required to be specified.

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